



A New Computation Technique of Glauert's Integrals

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Abstract—This paper deals with the study of particular functions of a complex variable which enable the easy calculation of Glauert's integrals.

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In aerodynamics, in order to apply the “Thin Airfoils Theory” by Glauert, the so-called Glauert's integrals are to be solved:

$$G_n = \int_0^\pi \frac{\cos(n\theta')}{\cos\theta' - \cos\theta} d\theta' = \pi \frac{\sin(n\theta)}{\sin\theta}, \quad (1)$$

$$G_n^* = \int_0^\pi \frac{\sin(n\theta') \sin\theta'}{\cos\theta' - \cos\theta} d\theta' = -\pi \cos(n\theta), \quad (2)$$

with $0 \leq \theta \leq \pi$. Glauert's integral (1) can be solved, for example, integrating in the complex plane the function

$$f(z) = \frac{z^n}{z^2 - 2z \cos\theta + 1}$$

with respect to z around the circle of unit radius $z = e^{i\theta'}$ and then equating the imaginary parts [1]. The integral (2) can be reduced to the form (1) by putting $\sin(n\theta') \sin\theta' = (1/2)[\cos(n-1)\theta' - \cos(n+1)\theta']$.

The solutions of Glauert's integrals can also be easily obtained, almost as a by-product, applying a new method that will be shortly described here.

Suppose we have on the strip $[(-c/2), (c/2)]$ on the x -axis a two-dimensional source distribution of strength $q(x)$ per unit length (c could be the chord of a symmetric airfoil). With $z = x + iy$, the complex velocity at any point of the complex plane is given by

$$u - iv = \frac{1}{2\pi} \int_{-c/2}^{c/2} \frac{q(x')}{z - x'} dx'. \quad (3)$$

Let us consider the behaviour of the complex velocity $u_0 - iv_0$ on the x -axis. The real part $u_0(x)$ is even on the whole axis, i.e., it does not change its sign through the axis, and vanishes at infinite distance from the strip

$$u_0(x) = \frac{1}{2\pi} \int_{-c/2}^{c/2} \frac{q(x')}{x - x'} dx'. \quad (4)$$

The imaginary part $-v_0(x)$, equal to zero outside the strip $[(-c/2), (c/2)]$, is odd on the strip, i.e., its sign changes through the strip, which must be treated as a double-faced strip

$$v_0(x) = \pm \frac{q(x)}{2}, \quad (5)$$

where the signs $+$ and $-$ are valid, respectively, on the upper and the lower face of the strip.

The properties exhibited by the velocity components on the x -axis due to a source distribution on a strip of the same axis are a particular case of a more general property of the functions of a complex variable.

Suppose that we want to find the function of a complex variable $F(z) = f + ig$ in the z -plane, assuming that the real part is even on the x -axis and equal to zero at infinity, and that the imaginary part is equal to zero on the x -axis except on the strip $[(-c/2), (c/2)]$, where it is odd and known. If $g_{0+}(x)$ is the value of the imaginary part on the upper surface of the double strip, according to equations (3)–(5) we obtain:

$$F(z) = f + ig = -\frac{1}{\pi} \int_{-c/2}^{c/2} \frac{g_{0+}(x')}{z - x'} dx', \quad (6)$$

$$f_0(x) = -\frac{1}{\pi} \int_{-c/2}^{c/2} \frac{g_{0+}(x')}{x - x'} dx'. \quad (7)$$

The above properties have been successfully applied to determine the complex velocity and the pressure distribution around a thin airfoil in unsteady incompressible flow [2].

To generalize this method to evaluate Glauert's integrals, now we can examine three particular functions which lead rapidly to the required solutions without any actual integration. Suppose that the variable $z = x + iy$ is obtained from the variable \bar{z} by means of the Zhukowsky transformation

$$z = \bar{z} + \frac{c^2}{16\bar{z}}. \quad (8)$$

If $\bar{z} = (c/4)e^{i\theta}$, equation (8) yields $z = (c/2)\cos\theta$. While \bar{z} varies on the circle of radius $c/4$, z varies on the double strip $[(-c/2), (c/2)]$; the upper face of the strip corresponds to $0 \leq \theta \leq \pi$ and the lower one to $-\pi \leq \theta \leq 0$.

On the double strip, we have

$$x = \frac{c}{2} \cos \theta; \quad dx = -\frac{c}{2} \sin \theta d\theta; \quad \int_{-c/2}^{c/2} \dots dx' = \int_0^\pi \dots \frac{c}{2} \sin \theta' d\theta'. \quad (9)$$

Consider first the function

$$F(z) = f + ig = \frac{c/2}{\sqrt{z^2 - c^2/4}}, \quad (10)$$

that satisfies the conditions for the validity of equations (6)–(7). On the double strip, the real part is zero and the imaginary part is an odd function

$$f_0 = 0; \quad g_{0+} = -\frac{1}{\sin \theta}. \quad (11)$$

Equations (6), (7), (9), (11) yield

$$F(z) = f + ig = \frac{1}{\pi} \int_0^\pi \frac{d\theta'}{2(z/c) - \cos \theta'}, \quad (12)$$

$$f_0 = 0 = \frac{1}{\pi} \int_0^\pi \frac{d\theta'}{\cos \theta - \cos \theta'}. \quad (13)$$

So, we have obtained the solution of integral (1) with $n = 0$: $G_0 = 0$.

Consider now the function

$$F(z) = f + ig = \frac{1}{((4/c)\bar{z})^n} = \cos(n\theta) - i \sin(n\theta). \quad (14)$$

On the double strip, we obtain

$$f_0 = \cos(n\theta); \quad g_{0+} = -\sin(n\theta). \quad (15)$$

Equations (6), (7), (9), (15) yield

$$F(z) = f + ig = \frac{1}{\pi} \int_0^\pi \frac{\sin(n\theta') \sin \theta'}{2(z/c) - \cos \theta'} d\theta', \quad (16)$$

$$f_0 = \cos(n\theta) = \frac{1}{\pi} \int_0^\pi \frac{\sin(n\theta') \sin \theta'}{\cos \theta - \cos \theta'} d\theta'. \quad (17)$$

Equation (17) gives the solution of integral (2) G_n^* .

Finally consider the function

$$F(z) = f + ig = \frac{1}{((4/c)\bar{z})^n} \frac{c/2}{\sqrt{z^2 - c^2/4}}. \quad (18)$$

Its value on the double strip is

$$[\cos(n\theta) - i \sin(n\theta)] \left(-\frac{i}{\sin \theta} \right),$$

and then we obtain

$$f_0 = -\frac{\sin(n\theta)}{\sin \theta}; \quad g_{0+} = -\frac{\cos(n\theta)}{\sin \theta}. \quad (19)$$

Equations (6), (7), (9), (19) give

$$F(z) = f + ig = \frac{1}{\pi} \int_0^\pi \frac{\cos(n\theta')}{2(z/c) - \cos \theta'} d\theta', \quad (20)$$

$$f_0 = -\frac{\sin(n\theta)}{\sin \theta} = \frac{1}{\pi} \int_0^\pi \frac{\cos(n\theta')}{\cos \theta - \cos \theta'} d\theta'. \quad (21)$$

From equation (21), we obtain the solution of integral (1) G_n .

REFERENCES

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